

QM Tipps 9

9.1 $\underline{x} = (x_1, \dots, x_N) = (q_1 \rightarrow q_N, p_1 \rightarrow p_N) \in \Gamma$

Kanonisches Ensemble $\left\{ \begin{aligned} w(\underline{x}) d\underline{x} &= \frac{1}{z(\beta)} e^{-\beta H(\underline{x})} \\ z(\beta) &= \int_{\Gamma} d\underline{x} e^{-\beta H(\underline{x})} \end{aligned} \right.$

$f: \Gamma \rightarrow \mathbb{R}$ "Observable", z.B. $H(\underline{x})$

$\langle f \rangle_{\beta} = \int_{\Gamma} f(\underline{x}) w(\underline{x}) d\underline{x} = \int \frac{f(\underline{x}) e^{-\beta H(\underline{x})}}{z(\beta)} d\underline{x}$

$S = -k \int w(\underline{x}) \log w(\underline{x}) d\underline{x} = k\beta \langle H \rangle_{\beta} + k \log z(\beta)$

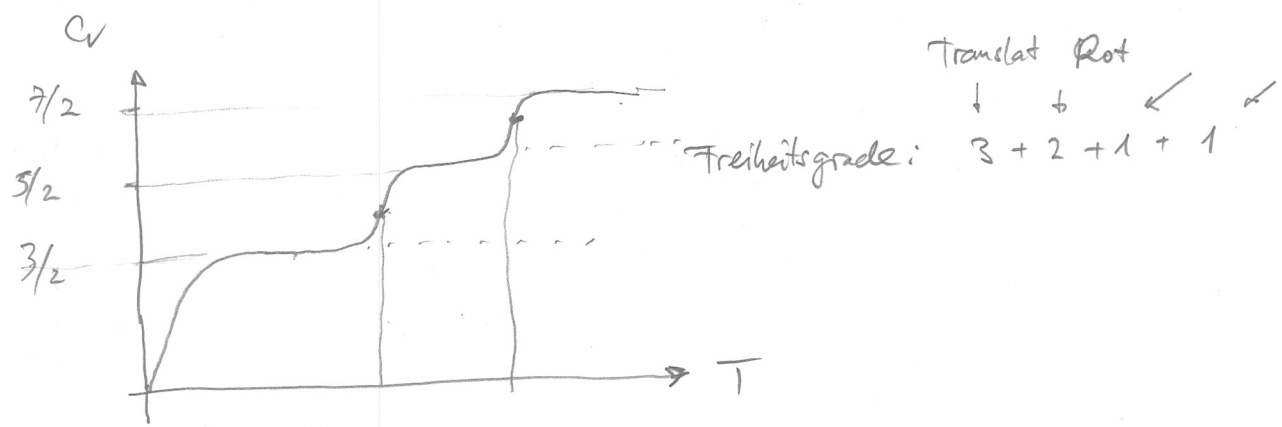
$F(\beta) = U - TS = -\frac{1}{\beta} \log z(\beta)$

$U = \langle H \rangle_{\beta} = -\frac{\partial}{\partial \beta} \log z(\beta) \quad c_v = \frac{dU}{dT} = \dots = -\frac{1}{kT^2} \frac{\partial U}{\partial \beta}$

Quantenstatistik

Sei $A = A^*$ Operator auf $\mathcal{H} = L^2(\mathbb{R}^{3N}, d^{3N}x)$

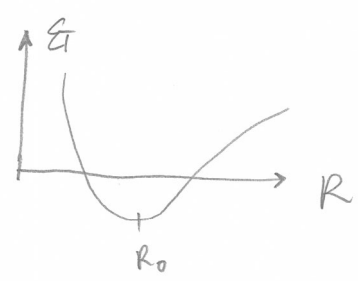
$\langle A \rangle = \text{tr}(A P_{\beta})$, $P_{\beta} = \frac{1}{z(\beta)} e^{-\beta H}$, $z(\beta) := \text{tr} e^{-\beta H}$
 $= \sum_n \langle \psi_n, e^{-\beta H} \psi_n \rangle$ $\{\psi_n\}$ ONB
 notwendig: H hat Punktspektrum



Ansatz: $C_v = \frac{3}{2}k + C_v^{vib} + C_v^{rot}$

$H^{rot} = \frac{I^2 \vec{L}^2}{2\Theta}$, $\Theta = M_{red} R_0^2 = \mu R_0^2$

$E_l = \frac{\hbar^2 l(l+1)}{2\Theta}$



$\text{tr } e^{-\beta H^{rot}} = \sum_{\alpha} \langle \varphi_{\alpha}, e^{-\beta H^{rot}} \varphi_{\alpha} \rangle$

Eigenfkt's: $m \cdot e^{im}$

$= \sum_{\alpha} e^{-\beta \frac{E(\alpha)}{2\Theta}} \langle \varphi_{\alpha}, \varphi_{\alpha} \rangle$ $\alpha = (l, m)$

$1 \stackrel{\text{normiert}}{=} \sum_{l,m} e^{-\beta \frac{l(l+1)}{2\Theta}} = \sum_{l=0}^{\infty} (2l+1) e^{-\frac{\beta l(l+1)}{2\Theta}}$
 analytisch nicht berechenbar
 Entartung in m

Betrachte Grenzwerte $\beta \rightarrow 0$, $\beta \rightarrow \infty$

$\sum_{l=0}^{\infty} f(l) \approx \int_0^{\infty} f(x) dx + \frac{1}{2} f(0) - \frac{1}{12} f'(0) + \dots$ Euler-Maclaurin

$(f(0) - f(\infty)) - \frac{1}{12} (f'(0) - f'(\infty))$

$$H(\phi_{\text{ext}}, \vec{A}_{\text{ext}}) = \frac{1}{2m} (\vec{p} - e \vec{A}_{\text{ext}})^2 - \frac{e^2}{|\mathbf{x}|} + e \phi_{\text{ext}}$$

Dipol approx: $E(\mathbf{x}, t) = E(t)$, kein B-Feld

$$H_K(t) = \frac{1}{2m} (\vec{p} - e \vec{A}(t))^2 - \frac{e^2}{|\mathbf{x}|}, \quad A(t) = - \int_{-\infty}^t \vec{E}(t') dt',$$

$$-A(t) = \vec{E}(t)$$

$$H_K(t) = -\frac{1}{2m} A - \frac{e^2}{|\mathbf{x}|} - e \underbrace{E(t) \cdot \mathbf{x}}_{\phi(t, \mathbf{x})}, \quad \vec{A} = 0$$

Eichtransformation: $\Lambda(\mathbf{x}, t)$

$$\vec{A} \rightarrow \vec{\tilde{A}} = \vec{A} + \vec{\nabla} \Lambda$$

$$\phi \rightarrow \tilde{\phi} = \phi - \partial_t \Lambda$$

$$\psi \rightarrow \tilde{\psi} = e^{+ie\Lambda} \psi$$

$$H_K: U_K(t, s) \text{ s.d. } \psi(t) = U_K(t, s) \psi(s)$$

$$H_R: U_K(t, s) \text{ s.d. } \tilde{\psi}(t) = U_K(t, s) \tilde{\psi}(s)$$

z.B. $U_K(t, s) = e^{ie\Lambda(\mathbf{x}, t)} U_K e^{-ie\Lambda(\mathbf{x}, s)}, \quad \Lambda(\mathbf{x}, t) = A(t) \cdot \mathbf{x}$

9.3

$$U_R(t,s) = S\text{-lim}_{N \rightarrow \infty} \prod_{j=1}^{N-1} e^{-i\varepsilon H_R(t-j\varepsilon)} \quad , \quad \varepsilon = \frac{t-s}{N}$$



$$\underbrace{e^{ie\Lambda(x,t)} e^{-i\varepsilon H_R(t)} e^{-ie\Lambda(x,t)} e^{ie\Lambda(x,t)} e^{-ie\Lambda(x,t-\varepsilon)} e^{ie\Lambda(x,t-\varepsilon)}}_{x e^{-i\varepsilon H_R(t-\varepsilon)}}$$

$$e^{ie\Lambda(x,t)-ie\Lambda(x,t-\varepsilon)}$$

$$e^{-i\varepsilon H_R(t)} \simeq 1 - i\varepsilon H_R(t) + \mathcal{O}(\varepsilon^2)$$

$$= e^{ie\varepsilon \Lambda(t)}$$

QM 1, Aufgabe 4.1

$$\left[e^{-ie\Lambda(x,t)} (\vec{p} - e\vec{A})^2 \psi = (\vec{p} - e\vec{A} + e\vec{\nabla}\Lambda)^2 e^{-ie\Lambda(x,t)} \psi \right]$$

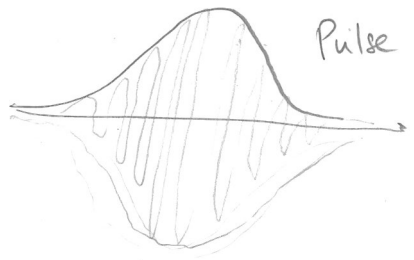
$$\underbrace{(\vec{p} - e\vec{A})^2}_{\parallel (\vec{p} - e\vec{A})} \psi = (\vec{p} - e\vec{A} + e\vec{\nabla}\Lambda)^2 \tilde{\psi}$$

$$\Rightarrow e^{-i\varepsilon \left[-\frac{1}{2m} (\vec{p} - e\vec{A})^2 - e\vec{E}(t)\vec{x} - \frac{e^2}{|\vec{x}|} \right]} e^{ie\varepsilon E(t)\vec{x}}$$

Birk: $e^A e^B = e^{\underbrace{A+B}_{\mathcal{O}(\varepsilon^2)}}$

$H(t) = H_0 + V(t) \quad V(t) = -e E(t) \vec{x}$

$E(t) = E_0 e^{-\frac{t^2}{\tau^2}} \begin{pmatrix} 0 \\ 0 \\ \cos \omega t \end{pmatrix}$



Zeit $s \rightarrow -\infty \quad \psi_{nem} \equiv \psi_i$
 $t \rightarrow +\infty \quad \psi_{n'el'n'} \equiv \psi_f$

$P_{fi}^{(2)}(t, s)$

(9.39) $P_{fi}^{(2)}(t, s) = \frac{1}{t^2} \left| \int_s^t dt' e^{i\omega_{fi} t'} \langle \psi_f, V(t') \psi_i \rangle \right|^2$
 $\omega_{fi} = \frac{E_f - E_i}{\hbar} \quad \sim \hat{z}$

ii) $P_{fi}^{(2)}(\infty, -\infty)$ mit $\psi_i = \psi_{100}$
 $\psi_f = \psi_{200} \quad \psi_{100} = \frac{1}{\sqrt{\pi a^3}} e^{-\frac{r}{a}}$

$\langle \psi_{100}, \hat{z} \psi_{200} \rangle = \frac{2^8}{3^3 \sqrt{2}} \cdot a \quad \int_{-\infty}^{\infty} dx e^{i\vec{k}\vec{x}} e^{-x^2/c^2} = 2 \frac{\sqrt{\pi}}{c} e^{-\frac{k^2 c^2}{4}}$

Vereinfache:

Einlaufendes Teilchen: $e^{-i\frac{\vec{p}}{\hbar}\vec{x}}$ $\vec{E} = E_0 \cos(\omega t) \hat{z}$

$P_{fi}^{(2)}(\infty, -\infty) = \langle e^{-i\frac{\vec{p}}{\hbar}\vec{x}}, \hat{V}(\omega_{fi}) e^{-\frac{r}{a}} \rangle \cdot \text{const}$
 $\vec{k} = \frac{\vec{p}}{\hbar} \quad \langle e^{i\vec{k}\vec{x}}, \hat{z} e^{-\frac{r}{a}} \rangle$

