

# 3D organisation of cells in pseudostratified epithelia

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**Pseudostratified epithelia have smooth apical and basal surfaces, yet along the apical-basal axis, cells assume highly irregular shapes, which we introduce as punakoids. They interact dynamically with many more cells than visible at the surface. Here, we review a recently developed new perspective on epithelial cell organisation. Seemingly random at first sight, the cell packing configurations along the entire apical-basal axis follow fundamental geometrical relationships, which minimise the lateral cell-cell contact energy for a given cross-sectional cell area variability. The complex 3D cell neighbour relationships in pseudostratified epithelia thus emerge from a simple physical principle. This paves the way for the development of data-driven 3D simulation frameworks that will be invaluable in the simulation of epithelial dynamics in development and disease.**

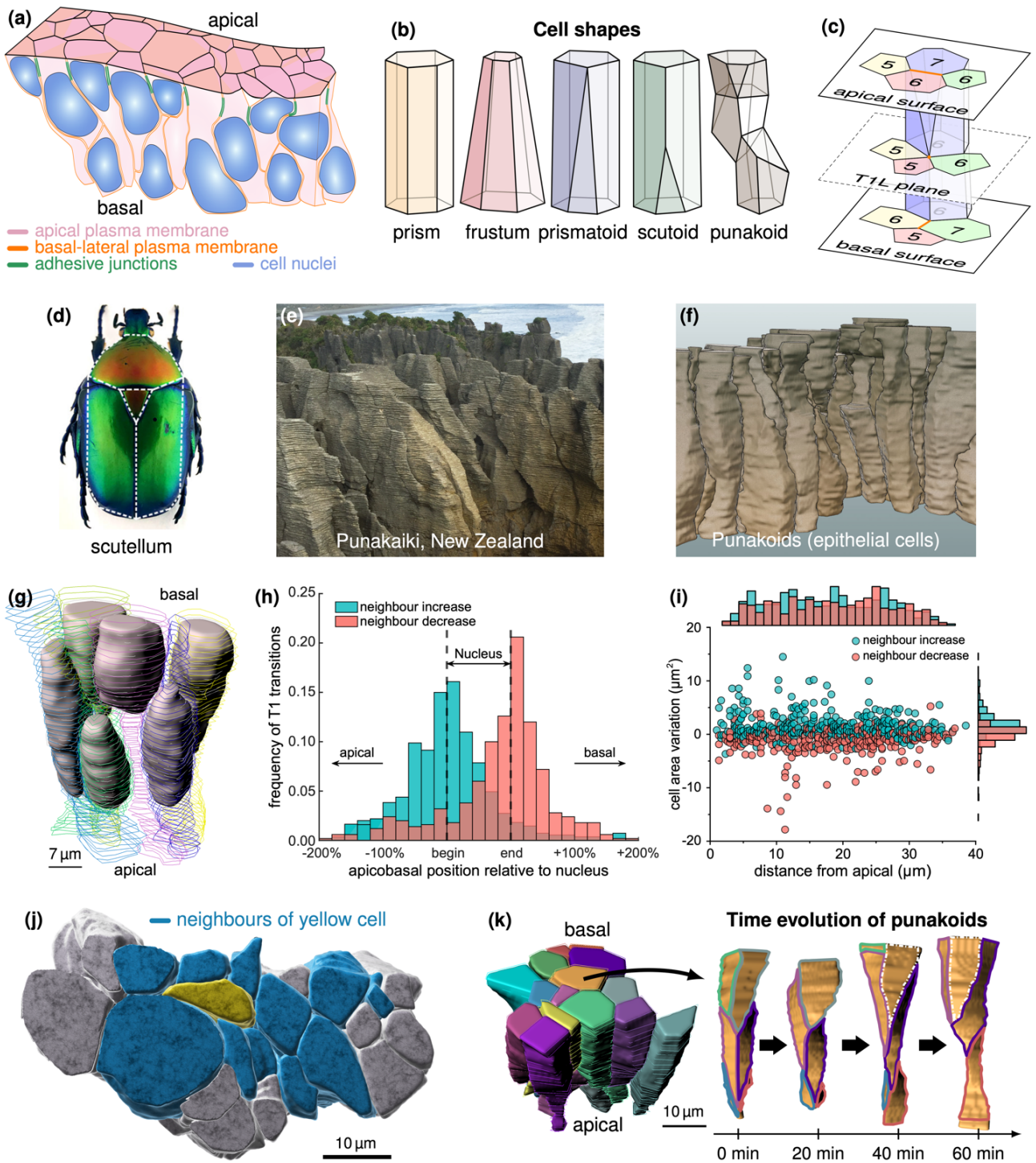
## Introduction

Epithelia are common to all animals and plants, and play a key structural role in tissue morphogenesis and the development of organ shapes. With more than 90% of cancers being of epithelial origin [1], there is an urgent need to uncover the principles of epithelial organisation and understand the basis for epithelial integrity and homeostasis. Epithelia achieve their structural function via their polarity (Figure 1a). On the outward-facing apical side, cells form a virtually impermeable barrier via a cadherin-based adhesion belt and tight junctions, while, on their basal side, they bind tightly to the basal lamina, a thin sheet composed of extracellular matrix (ECM) proteins [2-5]. Additional cell-cell junction complexes along the lateral sides provide further mechanical stabilisation. Recent advances in imaging provide insight into the physical principles according to which cell connectivity is organised in epithelia, and how it changes during morphogenesis and concomitant cell shape transitions.

## 3D epithelial cell shapes

Since the advent of light microscopy, epithelial surfaces have been studied in great detail, and this has revealed tight cell packing in polygonal lattices along the entire apical-basal axis [6-19]. As 3D segmentation of cells has become possible only very recently [17,18,20-25], 3D cell shapes have long

been depicted as prisms, which retain the same size and neighbour relationship along the entire apical-basal axis (Figure 1b). Cells in curved epithelial monolayers are commonly pictured as frustra (also termed bottle cells) as the apical and basal areas must differ. Differences in neighbour arrangements between the apical and basal side point to neighbour changes along the apical-basal axis in a range of epithelia [26]. Prismatoids accommodate the neighbour change at either surface. If the neighbour relationships change somewhere in between (Figure 1c), the cell shape is reminiscent of that formed by beetle scutum, scutellum and wings (Figure 1d), which led to the new term *scutoid* [15]. With up to 14 neighbour changes along the apical-basal axis [18], pseudostratified epithelial cells in developing mouse lungs, however, resemble more the pancake rock formations in Punakaiki at New Zealand's west coast (Figure 1e) than the back of beetles (Figure 1d). This novel complex geometry may therefore better be referred to as *punakoid* (Figure 1f). The defining characteristics of punakoids are 1) a quasi-polygonal surface, 2) a well-defined cell axis (in the pseudostratified epithelia reviewed here, the apical-basal axis), and 3) multiple neighbour changes along the axis. Like scutoids, the faces of punakoids are not necessarily planar, the edges not necessarily straight, and the entire shape not necessarily convex. So far, punakoids were found in the embryonic mouse lung, and their prevalence in other pseudostratified epithelia, although likely, remains to be demonstrated.



**Figure 1. 3D epithelial cell shapes.** (a) Pseudostratified epithelium with cell boundaries wrapping around wider nuclei. (b) Schematic depiction of cell shapes. (c) Cells alter their neighbour arrangements via T1L processes along the apical-basal axis. In a T1L process, two vertices that share an edge (orange) merge and decompose in a different direction such that neighbour relationships change. (d) The term scutoid was coined based on the shape of beetle scutum, scutellum and wings. (e) Pancake rocks at the beach of Punakaiki, NZ. (f) 3D segmented cells in the developing pseudostratified mouse lung epithelium (E12.5) resemble the Punakaiki rocks. (g) 3D shapes of six cells and their nuclei in an E12.5 pseudostratified mouse lung epithelium. (h) Neighbour relationships change predominantly at the apical and basal limits of the nuclei. (i) The lateral T1L processes are largely uniformly distributed along the apical-basal axis. Neighbour numbers tend to increase as cross-sectional cell area variation increases, and vice versa. (j) 3D cell neighbourhood extends further than apparent on the surface. (k) Time evolution of the contact areas between the central cell (orange) and its neighbours in a patch of 15 epithelial cells (left) over 60 minutes of explant culture. Panels a,g-k reproduced with modifications from [18], panel c from [15], published under the Creative Commons Attribution Non-Commercial 4.0 International License (CC BY-NC 4.0; <https://creativecommons.org/licenses/by-nc/4.0/>). Further reproduction of these panels would need to comply with the terms of this license.

As characteristic for pseudostratified epithelia [27], nuclei are found distributed along the entire apical-basal axis (Figure 1g) [18]. Neighbour changes occur mainly at the limits of the nuclei (Figure 1h), where the cross-sectional area changes the most [18]. An increase in the cross-sectional area correlates with a neighbour increase, and vice versa (Figure 1i). If the number of cells remains the same, neighbour relationships change via so-called T1 processes [28] (also referred to as rosette formation if more than four cells are involved [29]). We will refer to a neighbour transition along the apical-basal axis as lateral T1 process (T1L for short) (Figure 1c). The potentially large number of neighbour intercalations along their long axis lets the cells be in physical contact with others that, on the apical or basal surface, appear to be several cell diameters apart (Figure 1j). Cell-cell signalling can thus spread further than previously anticipated, and cells can read and average morphogen gradients over distances that were previously expected to require cell protrusions [30-33]. Much as on the apical surface [34], cell neighbour relationships further change dynamically over time along the entire apical-basal axis [17,18] (Figure 1k), thereby further increasing the distance over which signals can be sensed, exchanged, and averaged. But what leads to these unexpectedly complex and dynamic 3D cell shapes, and what determines cell neighbour relationships?

### Surface area minimisation

Epithelial cells are often compared to soap bubbles. Soap bubbles famously minimise their surface area and assume a spherical shape in isolation. Motivated by the tight packing of soap bubbles in foams, there has been a long-standing interest in optimal packing solutions that minimise the overall surface area. In the 19<sup>th</sup> century, Lord Kelvin proposed that tetrakaidekahedra minimise the overall surface area if all soap bubbles have the same volume. Flattened 14-sided tetrakaidekahedra with hexagonal apical and basal surfaces are found in the multi-layered, stratified epidermis of the skin [22]. However, a more efficient packing of equally sized cells has since been described [35], and the complex shapes of cells in single-layered pseudostratified epithelia (Figure 2a) certainly do not minimise the overall surface area for the given cell volume. What then governs their shape and neighbour relationships?

### Striking regularities in cell arrangements

The 3D cell neighbour relationships can be understood by considering single 2D planes, perpendicular to the apical-basal cell axis [18]. In the following, we will therefore first discuss the neighbour relationships in these 2D planes (Figure 2).

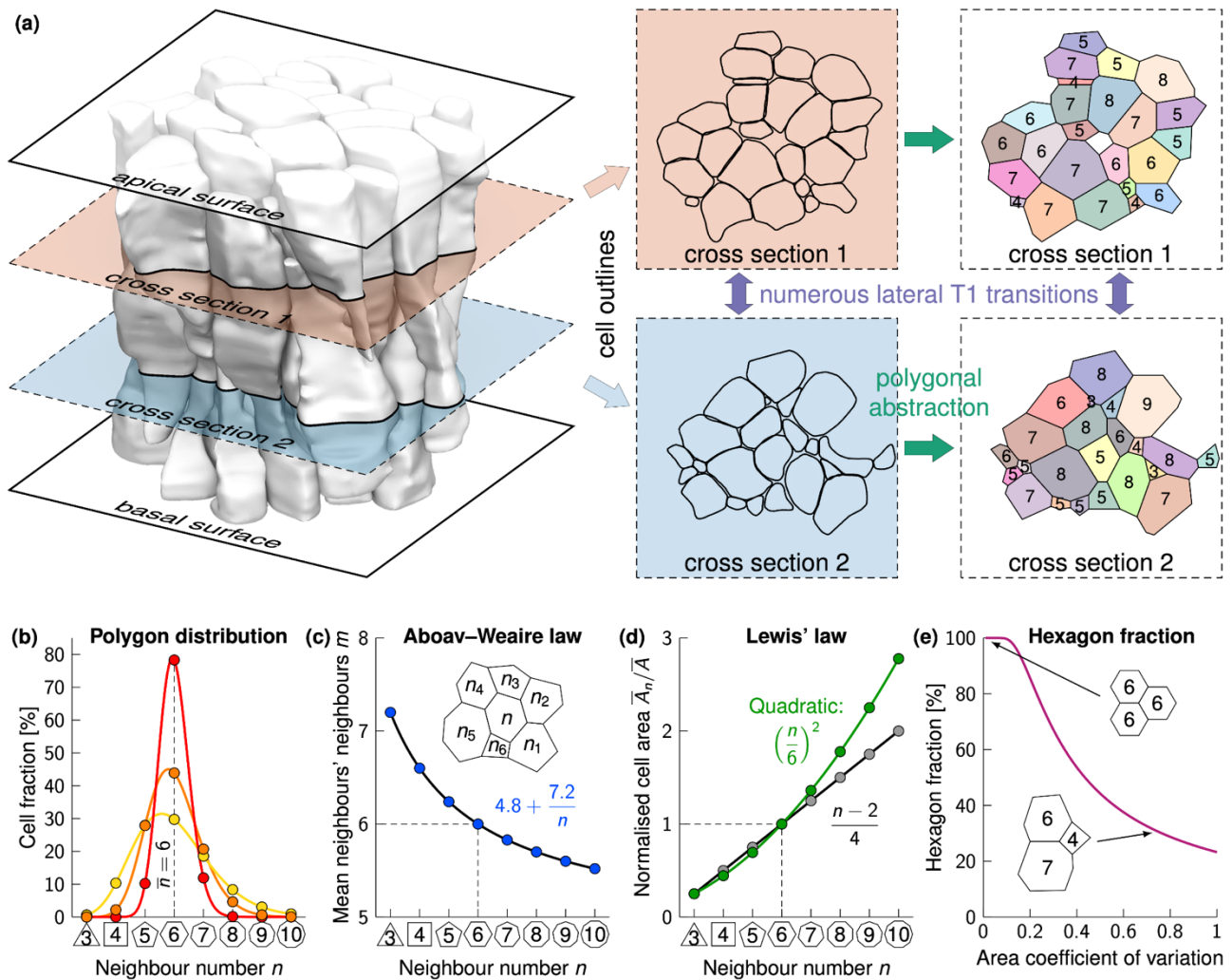
Striking regularities that have long been known and that are found on *all* apical and basal surfaces studied to date [6-16,19], have recently been reported also in all planes along the apical-basal axis (Figure 2a), even though neighbour relationships change between individual cells [18]. First and foremost, cells in any 2D plane have on average (close to) six neighbours, albeit the neighbour number distributions (Figure 2b) differ significantly among epithelia and between planes. This can be accounted to topological constraints in 2D contiguous polygonal lattices, and follows directly from Euler's polyhedron formula [10,36]. If three edges meet at each junction, the mean neighbour number in infinite lattices is exactly six,  $\bar{n} = 6$ . The average declines as the number of edges per vertex increases, to  $\bar{n} = 4$  if four edges meet in each junction. Locally, the average deviates from six in epithelia (Figure 2c) and follows a phenomenological relationship known as Aboav-Weaire's law [37],

$$m = \frac{1}{n} \sum_{i=1}^n n_i = a + \frac{b}{n} \quad (1)$$

which relates the number of neighbours,  $n$ , of a central cell to the average one of its neighbours,  $m$  (Figure 3c, inset). In epithelia, the parameter values fall into the range  $a \in [4.5, 5.5]$  and  $b \in [4.5, 9.5]$  [38]. Finally, the relative average apical area,  $\bar{A}_n$ , of cells with  $n$  neighbours with respect to the average area of all cells,  $\bar{A}$ , linearly increases with  $n$  (Figure 2c, black line), a phenomenological relation termed Lewis' law [6],

$$\frac{\bar{A}_n}{\bar{A}} \approx \frac{n-2}{4}. \quad (2)$$

Initially, Lewis' law has been accounted to entropy maximisation [36], but this has subsequently been ruled out [39,40]. Many other hypotheses have been explored to explain epithelial organisation. According to topological arguments, sequential cell division results in the observed frequencies of neighbour numbers (Figure 2b) [10,41]. However, this argument does not explain the emergence of cells with less than 5 neighbours, and predicts proliferative epithelial tissues to have about 45% hexagons. The hexagon frequencies, however,



**Figure 2. Phenomenological laws in epithelial cell arrangements.** (a) Slicing through a mouse embryonic lung epithelium (shaded cells [18]) repeatedly along the apical-basal axis reveals the complex shape and packing structure of punakoids. Even over short distances, numerous T1L transitions occur, leading to vastly different cell neighbourships. Colours in the last column indicate cell identity, values are cross-sectional neighbour numbers. (b) The average neighbour number,  $\bar{n}$ , is (close to) 6 in all epithelia, even though the hexagon fraction and the neighbour number distribution vary. (c) All quantified epithelial tissues follow Aboav-Weaire's law. (d) All quantified epithelial tissues follow either Lewis' law (grey line), or the quadratic law (green line), or lie in between. (e) The hexagon frequency declines for increasing area variability according to the magenta curve.

decrease with increasing variability in the cell cross-sectional areas (Figure 2e), and reported values range from 30% to 80% [14]. Contrary to the assumptions of the topological model, cells rearrange their boundaries until they reach a mechanical equilibrium [9]. By altering the relative cell-cell adhesion strength and cortical tension, the full range of neighbour relationships can be reproduced in vertex models and similar model setups [9,14,41-46]. In a small subset of the parameter space, Lewis' law emerges [9,43,44]. Lewis' law and the entire range of measured neighbour frequencies can be reproduced also using Voronoi tessellations, but again only for the subset of the parameter space that yields the right level of tessellation irregularity [13,47]. So, why do all epithelia follow those two phenomenological laws?

### Minimisation of the lateral cell-cell contact energy determines cell neighbour relationships

As cells reach the mechanical equilibrium quickly (in less than a minute [9]), the polygonal lattices that one observes when cutting the epithelium in any plane (Figure 2a) represent a mechanical equilibrium, i.e., a state of minimal energy. At first sight, the highly irregular shapes of epithelial cells may appear inconsistent with surface energy minimisation, as observed in foam. However, by following Lewis' law and Aboav-Weaire's law, epithelial cells still minimise the lateral surface area for the given irregular cell volume distribution [14,38]. Thus, in each plane along the apical basal-axis (Figure 2a), cells minimise the total perimeter for the enclosed cross-sectional areas.

As regular polygons have the smallest perimeter per enclosed area, a lattice composed of regular polygons will have the smallest total perimeter. If all cells had the same cross-sectional area, a regular hexagonal lattice would be most favourable. However, cellular processes constantly alter the cross-sectional areas, and the combined cell-cell contact surface energy is lower with mixed cross-sectional cell areas [14]. Even the hexagonal ommatidia in the *Drosophila* eye are each composed of 21 differently-sized apical cell areas, which are predominantly not hexagonal [48]. The arrangement into hexagonal ommatidia relies on the careful adjustment of cell adhesion, cortical tension, and cell dilation [49-51]. As mixed cross-sectional cell areas are most favourable, epithelial cells easily disperse from a clone with smaller cells, while they remain clustered without such a cell size difference relative to the surrounding tissue, potentially facilitating the spreading of tumour cells [16].

For the distribution of cross-sectional cell areas found in epithelial tissues, perfectly regular contiguous lattices cannot form. By following Aboav-Weaire's law, the internal angles of the polygons are closest to those of a regular polygon while still adding up to  $360^\circ$  at each junction [38]. By following Lewis' law, the side lengths are most similar [14]. Equal side lengths are obtained if the cross-sectional cell areas follow a quadratic relation (Figure 3d, green line) of the form

$$\frac{\overline{A_n}}{\overline{A}} \approx \frac{n \tan(\pi/6)}{6 \tan(\pi/n)} \approx \left(\frac{n}{6}\right)^2 \quad (3)$$

The quadratic relationship, however, emerges only at a high area variability, as found on the apical side of embryonic lung tubes [14,18]. Finally, a novel relationship that all epithelia follow emerges from the drive to the most regular polygonal shape, and relates the fraction of hexagons to the apical area variability, measured by the coefficient of variation ( $CV = \text{std}/\text{mean}$ ) (Figure 2e) [14]. Interestingly, even puzzle cells in plants, which derive their name from their highly irregular shape, reminiscent of puzzle pieces, follow Lewis' law [52]. This is still consistent with a minimisation of the cell perimeter because the puzzle shape emerges in an effort to minimise stress in large cross-sections only after the cells have stopped dividing and the neighbour relationships have been fixed by the rigid cell walls [53]. In summary, the minimisation of lateral cell-cell contact energy defines the polygonal shape of each cell cross-section, and thus the cell neighbour relationships. Changes in the relative cross-sectional areas along the apical-basal axis or over time drive cell neighbour changes [14,18]. But why

do the cross-sectional areas change – or differently put, what defines the 3D cell shape?

### Impact of tissue curvature on cell neighbour relationships

If the two principal curvatures of the tissue surface change differently along the apical-basal axis, such as in sufficiently thick epithelial tubes, then the cell aspect ratio changes along the apical-basal axis. To maintain a regular polygonal cross-sectional cell shape, neighbour relationships have to change. This curvature effect has been proposed to result in scutoid cell shapes in epithelial tubes [15]. Curvature-driven T1L processes should then on average occur more frequently, i.e., for lesser curvature fold-changes  $\kappa_2/\kappa_1$ , the higher the cell neighbour number,  $n$ , in local cross sections [18]:

$$\frac{\kappa_2}{\kappa_1}(n) = \left(1 - \frac{\alpha(2+\cos\alpha)}{n \sin\alpha(1+2\cos\alpha)}\right)^{\pm\frac{\pi}{2}}, \quad \alpha = \frac{2\pi}{n}. \quad (4)$$

In the tubular embryonic mouse lung epithelium, no such systematic  $n$ -dependency is observed [18]. Moreover, in planar monolayers and in spherically shaped epithelia, where the principal curvatures change equally, T1L transitions are nonetheless still found [15,18,19]. Effects other than tissue curvature must thus dominate in these epithelia.

### Determinants of 3D cell shape and neighbour relationships

The shape of cells in single-layered epithelia can range from a cuboidal to highly elongated columns with large aspect ratio, and apical or basal constriction can further affect the cell shape [20,21,54]. In highly elongated cells, the diameter of a spherical nucleus would be larger than the diameter of a cylindrical cell. Accordingly, both the cells and the nuclei deform [55]. The cell is wider where the nucleus is present, and the remaining part of the cell is necessarily much thinner (Figure 1g). At the apical and basal limits of the nucleus, there is a sharp change in the cell cross-sectional areas, and most changes in neighbour relationships are found in this transition zone (Figure 1h). Epithelia with an average cell diameter smaller than the maximal nuclear diameter can thus be expected to have many more neighbour changes than those with wider cells.

Given the narrow columnar shape, there is insufficient space to accommodate all nuclei simultaneously in the same plane. Accordingly, the

nuclei of neighbouring cells are found at different positions along the apical-basal axis (Figures 1a, g) [18], a configuration referred to as pseudostratification [27]. As mitosis is restricted to the apical side [27,56], nuclei actively move towards the apical side during the G2 phase, and are pushed towards the basal side as the cell exits mitosis, in a process called interkinetic nuclear migration (IKNM) [55,57,58]. As the nuclei translocate between the apical and basal side during the cell cycle, the cell cross-sectional areas and connectivities continuously change (Figure 1k). An increase in the cross-sectional area increases the chance of a neighbour increase and vice versa (Figure 1i). Neighbour changes are less frequent close to the basal surface of tube segments, where cells remain wider throughout the cell cycle, but are otherwise uniformly distributed along the apical-basal axis [18]. Consistent with a stochastic basis to the 3D organisation of epithelial cells, the number of T1L per cell is Poisson-distributed [18].

But why would epithelial cells adopt such an elongated cell shape? Independent of the increased number of dynamically changing cell contacts, a smaller cell diameter can increase the precision of morphogen-based patterning [30]. Interestingly, several diffusible morphogens and growth factors, including Fibroblastic Growth Factor (FGF), Sonic Hedgehog (SHH), Bone Morphogenetic Protein (BMP) / transforming growth factor-beta (TGF- $\beta$ ), and WNT, have been observed to affect cell height, presumably via an effect on cell tension and/or cell-cell adhesion [59-65].

### **Discussion: Towards 3D cell-based tissue simulations of epithelial dynamics**

As the complex, dynamic 3D organisation of cells in growing epithelia is governed by simple physical concepts, computer simulations present powerful tools to understand the emergent properties of epithelia [66], including IKNM and its effects [67-71]. Cellular Potts models, which represent a generalisation of the Ising model to cells, have long been used to simulate complex 3D cell shapes [72-74]. Vertex models have been developed to specifically represent epithelia in 3D, but without resolving the complex irregular shapes of epithelial cells [64,75-77]. In 2D, several vertex-based models with higher cell boundary resolution have been developed to enable more complex cell shapes [78-82], and to represent individual cell boundaries and the interstitial volume [83-86]. A recent hybrid version between a spheroid and a vertex model allows for a 3D vertex model with an intermediate

vertex that enables a neighbour transition along the apical-basal axis [87]. To make full use of the available 3D imaging data, efficient, high-resolution vertex-based simulation frameworks are now required. A first such simulation framework that represents cells by individual, deformable meshes has recently been developed [88,89]. In combination with quantitative 3D imaging data, this now paves the way to a more detailed understanding of epithelial cell dynamics in development and disease. Vice versa, cell shape data can be used to infer force fields and to predict bias in cell division as cells divide perpendicular to the longest axis of their apical surface [41,90-94]. With such tools at hand, it may become feasible to address open questions regarding the maintenance and loss of epithelial integrity and cell polarity, for instance in tumour growth and mesenchymal-to-epithelial transitions.

### **Declaration of Interest**

The authors declare no conflict of interest.

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